



"Amplitude scattering matrix"

More practical approach: Stokes vector and the phase matrix

$$I \propto \left\langle E_{\parallel} E_{\parallel}^{*} + E_{\perp} E_{\perp}^{*} \right\rangle$$
$$Q \propto \left\langle E_{\parallel} E_{\parallel}^{*} - E_{\perp} E_{\perp}^{*} \right\rangle$$
$$U \propto \left\langle E_{\parallel} E_{\perp}^{*} + E_{\perp} E_{\parallel}^{*} \right\rangle$$
$$V \propto i \left\langle E_{\parallel} E_{\perp}^{*} - E_{\perp} E_{\parallel}^{*} \right\rangle$$





Phase matrix depends on particle size, shape, composition, orientation, wavelength of light, illumination geometry

Phase matrix is simplified in special cases:

Rayleigh scattering

$$\begin{array}{cccccc} P_{11} & P_{12} & 0 & 0 \\ P_{12} & P_{11} & 0 & 0 \\ 0 & 0 & P_{33} & 0 \\ 0 & 0 & 0 & P_{33} \end{array}$$

Mie scattering (= spheres)

$$\begin{array}{cccccc} P_{11} & P_{12} & 0 & 0 \\ P_{12} & P_{11} & 0 & 0 \\ 0 & 0 & P_{33} & P_{34} \\ 0 & 0 & -P_{34} & P_{33} \end{array}$$

More general case. Randomly oriented nonspherical particles, equal amount of particles and their mirror particles.

If incident light is *unpolarized* (*Q*, *U*, *V* = 0): Intensity of scattered light $I_s \propto P_{11}(\Theta)$ Degree of linear polarization $-\frac{P_{12}(\Theta)}{P_{11}(\Theta)}$

Rayleigh Scattering Geometry





Rayleigh Scattering formulae

$$Q_{s} = \frac{8}{3} x^{4} \left| \frac{m^{2} - 1}{m^{2} + 2} \right|^{2} \qquad \varpi \sim x^{3}$$
$$Q_{a} = 4x \operatorname{Im} \left(\frac{m^{2} - 1}{m^{2} + 2} \right)$$

Phase Matrix

$$\begin{bmatrix} 1 + \cos^2 \Theta & \cos^2 \Theta - 1 & 0 & 0 \\ \cos^2 \Theta - 1 & 1 + \cos^2 \Theta & 0 & 0 \\ 0 & 0 & 2\cos \Theta & 0 \\ 0 & 0 & 0 & 2\cos \Theta \end{bmatrix}$$

Cloud Water Absorption in the microwave



Mie solution; the result

$$Q_{\rm e} = rac{2}{x^2} \sum_{n=1}^{\infty} (2n+1) \Re(a_n+b_n) ,$$

$$Q_{\rm s} = rac{2}{x^2} \sum_{n=1}^{\infty} (2n+1)(|a_n|^2+|b_n|^2)$$
 ,

12.23)
$$S_1 = \sum_n \frac{2n+1}{n(n+1)} (a_n \pi_n + b_n \tau_n)$$

(12.24)
$$S_2 = \sum_n \frac{2n+1}{n(n+1)} (a_n \tau_n + b_n \pi_n)$$

Expansion coefficients of the scattered field: $a_n = \frac{m\psi_n(mx)\psi_n'(x) - \psi_n(x)\psi_n'(mx)}{m\psi_n(mx)\xi_n'(x) - \xi_n(x)\psi_n'(mx)}$ $b_n = \frac{\psi_n(mx)\psi_n'(x) - m\psi_n(x)\psi_n'(mx)}{\psi_n(mx)\xi_n'(x) - m\xi_n(x)\psi_n'(mx)}$

Here ψ , ξ are Riccati-Bessel functions and π , τ are functions of Legendre polynomials, x is the size parameter and m the complex refractive index.

The series expansions converge at large nand are, in practice, truncated at some $n \approx x \dots 2x$.

Mie solution; using the result

Typical Inputs:

x (size parameter)*m* (complex refractive index)

Typical Outputs:

 $S_1(\Theta)$, $S_2(\Theta)$: Amplitude scattering matrix (complex) Qe, Qs : Ext, Scat efficiencies

$$\sigma_{e} = Q_{e}\pi r^{2}$$

$$\sigma_{s} = Q_{s}\pi r^{2}$$

$$\rho = \frac{4}{3}\pi r^{3}\rho_{m}n$$

$$P_{11} = \frac{2}{Q_{s}x^{2}} \left(|S_{1}|^{2} + |S_{2}|^{2} \right)$$

$$P_{12} = \frac{2}{Q_{s}x^{2}} \left(|S_{2}|^{2} - |S_{1}|^{2} \right)$$

$$P_{33} = \frac{2}{Q_{s}x^{2}} \left(S_{2}^{*}S_{1} + S_{1}^{*}S_{2} \right)$$

$$P_{34} = \frac{2}{Q_{s}x^{2}} \operatorname{Im} \left(S_{2}^{*}S_{1} - S_{1}^{*}S_{2} \right)$$

Phase Function of water spheres (Mie theory)



High Asymmetry Parameter

Low Asymmetry Parameter

Mie Theory

• Exact Qs, Qa for spheres of some x, m.

$$Q_{\rm e} = \frac{2}{x^2} \sum_{n=1}^{\infty} (2n+1) \Re(a_n + b_n) , \qquad (12.23)$$

$$Q_{\rm s} = \frac{2}{x^2} \sum_{n=1}^{\infty} (2n+1)(|a_n|^2 + |b_n|^2) , \qquad (12.24)$$

- a, b coefficients are called "Mie Scattering coefficients", functions of x & m. Easy to program up.
- "bhmie" is a standard code to calculate Q-values in Mie theory.
- Need to keep approximately $x + 4x^{1/3} + 2$ terms for convergence



Mie Theory Results for ABSORBING SPHERES



Variations of SSA with wavelength

